

interpolated value of the phase space density at grid point (i, j) in phase space with the corresponding interpolated data, f_{data} , at the same point, and σ is the weighting associated with $\log(f_{data})$.

To handle the nonrandom scatter of the data, we have incorporated two "robust" techniques [Press et al., 1986], a truncation ceiling for the evaluation of single point χ^2 , and a constant σ . Even with a smoothed β spline interpolation with zero tension [Bartels et al., 1983], it is possible for the model to deviate so markedly from the data in certain regions that these regions dominate χ^2 . Thus we also allow restricting the fit domain in M - L space.

Reference Fits (1-4)

Since electric diffusion dominates below 1 keV/nT, and because the data above 1 keV/nT is also above $L=6$ where losses are too small to constrain the fit, we did not adjust the magnetic diffusion coefficient, setting it to nominal values, $D_{LL}^M = 2.3 \times 10^{-15} R_E^2/s$ (SSF). Thus we began by fitting the reference model, allowing only the amplitude of the electric diffusion coefficient, D_{LL}^E , to vary, fixing D_{LL}^M , and setting the domain to subsets of the entire data set as summarized in Table 3. Single point χ^2 values were truncated at 10 (a 3 order of magnitude deviation in log space). Fit 1 over H^+ , He^+ , and He^{++} is shown in Plate 1 with the data in the first column, the fit in the second column and the ratio of data to fit in the third column. The logarithmic scale extends from 10^{-1} (blue)

the data correspond to 1 and 300 keV/e constant energy boundaries. The black lines indicate the domain of the data used in the fit. The entire region is modeled (so that the ratio exists wherever the data exist), but modeled phase space contours for $f < 10^{-1}$ are not shown for the sake of clarity.

The data and the reference model agree in the $L > 5$, $M > 0.1$ keV/nT region. The upper boundary at $L=7.5$ agrees perfectly since the data were the boundary condition for the model. Since the qualitative behavior of H^+ is similar to that of both He^+ and He^{++} , we will discuss H^+ in detail and refer to helium only where it differs. The H^+ with $M > 1$ keV/nT diffuse to lower L shells with comparatively little loss since they are above the peak of the charge exchange cross section, $E \sim 60$ keV (Appendix B on microfiche). Eventually Coulomb drag sweeps the penetrating particles to lower energy below $L \sim 5$. The spectral peak seen just to the left of the black line at $M < 0.02$ keV/nT, $L < 4.5$ is not an artifact, but related to convection electric field transport.

Convective Transport

In the first column of Plate 2 we illustrate the convective transport of ions into the ring current using a dipole magnetic field and Volland-Stern [Volland, 1973; Stern, 1973] convection electric field model (including corotation) with a Maynard and Chen [1975] and

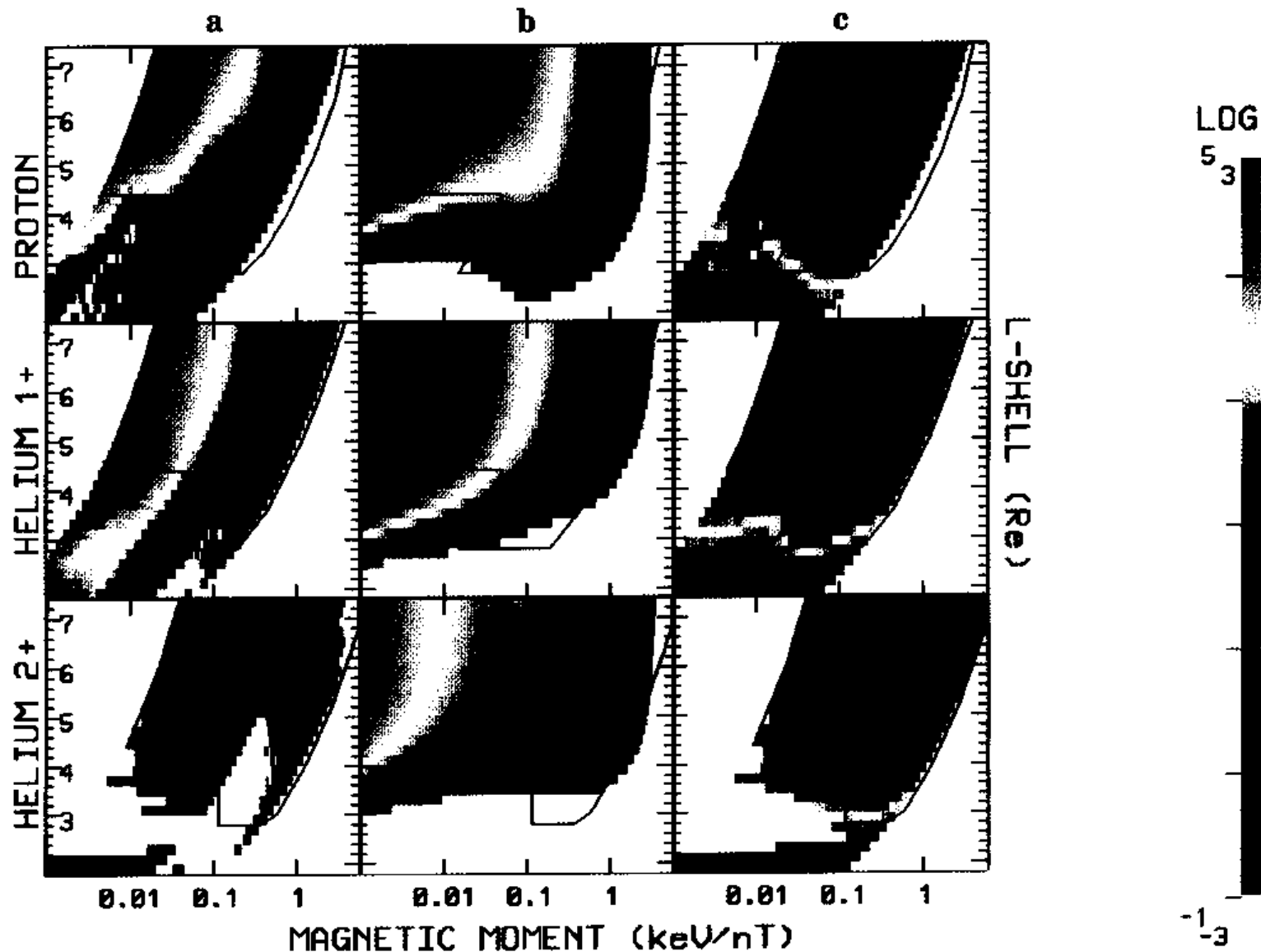


Plate 1. Measured phase space density in s^3/km^6 (column 1), reference model (column 2), and the ratio of data/reference (column 3). The range of intensities displayed is logarithmic from -1 (blue) to 5 (red) for columns 1 and 2, and from -3 to 3 for column 3.