



Figure 2. Logarithms of the phase space density ratio of data/model for fit 0, fit 1, and fit 2 in columns 1–3, respectively. The range of intensities displayed is logarithmic $10^{-1.5}$ (white)– $10^{1.5}$ (black) with contours every half-decade. Perfect agreement contour is shown dotted. Solid line outlines fitting region as in Figure 1.

The value for χ^2 should be taken as a relative measure of the “sharpness” of the χ^2 minimum, and not as confidence limits, both because χ^2 was evaluated with equal weights for each point and because deviations from the fit were not normally distributed. Nonetheless, the value of D^E is well constrained by the data, and also within the range of previous measurements [West *et al.*, 1981], though about three times larger than the value of $2.3 \times 10^{-15} R_E^2/s$ used by [Spjeldvik, 1977].

Fit 1: External Source With Stall Correction

We began with externally driven electric field fluctuations as before, $D^E \propto L^6$, but included the effect of the corotation $\vec{E} \times \vec{B}$ drift on the drift frequency, as well as our formulation for the electric diffusion coefficient (equation (18)). Since the electric field shielding has an expected time constant, $0.5 < \tau < 3$ hours, [Vasyliunas, 1972; Kaye and Kivelson, 1981; Anderson *et al.*, 1991], we assumed no additional radial dependence in the amplitudes, fixing q , L_0 and p , but fitting the spectral coefficients n , above and m , below Ω_E .

Thus we made a four parameter fit, holding other parameters at fit 0 values, and found a 25% improvement in χ^2 .

In Figure 1 the second and third columns show the modified and standard diffusion coefficients from fit 6 and reference models, respectively. The narrow feature in the second column follows the “stall” energy of the ions, where we have included only the corotation electric field and neglected the convection electric field contribution. In more realistic electric and magnetic field models, which produce radial drift, not included here, the ion drift velocity does not decrease to zero but merely passes through a minimum at this location. In both cases the drift frequency is reduced which results in less diffusion if there is reduced power at lower frequencies.

We plot the logarithm of the ratio of data to the model in column 2 of Figure 2. We see that the major effect of the new transport coefficient is to produce better agreement at lower L shells than seen in the reference model. In particular, we reproduce the “peninsula” of diffusing protons above 1 keV/nT very well. The over prediction of the model “white spot” near 1 keV/nT and $L \sim 5$ has improved only slightly.